

Thermal Pure Quantum States of Many-Particle Systems: *Supplemental Material*

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Practical formulas for numerical computation using the grand-canonical thermal pure quantum (gTPQ) states are presented.

In numerical computations, one has to introduce a certain cutoff to make $\dim \mathcal{H}$ finite. Then, $\hat{g} \equiv (\hat{H} - \mu\hat{N})/V$ has the maximum eigenvalue g_{\max} , and we can arbitrarily take a number l such that $l \geq g_{\max}$. Using l , we define

$$|k\rangle \equiv \sum_{\nu} z_{\nu} (l - \hat{g})^k |\nu\rangle \quad (k = 0, 1, 2, \dots), \quad (1)$$

and the normalized one, $|\psi_k\rangle \equiv (1/\sqrt{Q_k})|k\rangle$, where $Q_k \equiv \langle k|k\rangle$. If \hat{g} is replaced with \hat{H}/V , $|\psi_k\rangle$ reduces to the mTPQ state of Ref. [1]. Hence, we simply call $|k\rangle$ and $|\psi_k\rangle$ mTPQ states. We have

$$e^{V\beta l/2} |\beta\mu V\rangle = \sum_{k=0}^{\infty} \frac{(V\beta/2)^k}{k!} |k\rangle = \sum_{k=0}^{\infty} R_k |\psi_k\rangle, \quad (2)$$

where $R_k \equiv (V\beta/2)^k \sqrt{Q_k}/k!$. We can prove that this sum is uniformly convergent on any finite interval of β . We can also prove that R_k takes significant values only for k such that

$$\langle \psi_k | \hat{g} | \psi_k \rangle = \langle \hat{g} \rangle_{\beta\mu V}^{\text{TPQ}} + O(1/V). \quad (3)$$

As k moves away from such values, R_k vanishes exponentially fast. This means that $|\beta\mu V\rangle$ is almost composed of $|\psi_k\rangle$'s which represent the same equilibrium state as that represented by $|\beta\mu V\rangle$. Therefore, one can terminate the sum at a finite number k_{term} , which depends on the largest β of interest, β_{\max} . For any $\beta_{\max} = \Theta(1)$, we can show that $k_{\text{term}} = \Theta(N)$. Since $|\psi_1\rangle, |\psi_2\rangle, \dots, |\psi_{k_{\text{term}}}\rangle$ can be obtained iteratively by simply multiplying $(l - \hat{g})$ with a random vector k_{term} times [1], one can obtain $|\beta\mu V\rangle$ by multiplying $(l - \hat{g})$ repeatedly $\Theta(N)$ times.

We can obtain more practical formulas. Let

$$\{\hat{A}\}'_{\beta\mu V} \equiv \sum_{k=0}^{\infty} \left[\frac{(V\beta)^{2k}}{(2k)!} \langle k | \hat{A} | k \rangle + \frac{(V\beta)^{2k+1}}{(2k+1)!} \langle k | \hat{A} | k+1 \rangle \right]. \quad (4)$$

When $[\hat{A}, \hat{g}] = 0$, we have

$$\{\hat{A}\}'_{\beta\mu V} = \langle \beta\mu V | \hat{A} | \beta\mu V \rangle. \quad (5)$$

For example, $\{\hat{1}\}'_{\beta\mu V} = \langle \beta\mu V | \beta\mu V \rangle$. Even when $[\hat{A}, \hat{g}] \neq 0$, the above equality is almost exact. To see this, we take $x = \{\hat{A}\}'_{\beta\mu V} \equiv \{\hat{A}\}'_{\beta\mu V} / \{\hat{1}\}'_{\beta\mu V}$ and $y = \langle \hat{A} \rangle_{\beta\mu V}^{\text{ens}}$ in inequality (3) of the paper, and evaluate

$$E_V(A)^2 \equiv \overline{\{\hat{A}\}'_{\beta\mu V}^{\text{TPQ}} - \langle \hat{A} \rangle_{\beta\mu V}^{\text{ens}})^2} \quad (6)$$

as

$$E_V(A)^2 \leq \frac{\langle (\Delta\hat{A})^2 \rangle_{\beta\mu V}^{\text{ens}}}{\exp[V\beta\{j(0, \mu; V) - j(T, \mu; V)\}]}. \quad (7)$$

This shows that

$$E_V(A)^2 \leq V^{2m}/e^{\Theta(V)} \quad (8)$$

and $\{\hat{A}\}'_{\beta\mu V}^{\text{TPQ}} \xrightarrow{P} \langle \hat{A} \rangle_{\beta\mu V}^{\text{ens}}$ exponentially fast and uniformly. With this formula one needs only to calculate $\langle k | \hat{A} | k \rangle$ and $\langle k | \hat{A} | k+1 \rangle$ for all $k \leq k_{\text{term}}$ to obtain the results for *all* $\beta \leq \beta_{\max}$.

When computer resources are not sufficient to treat large enough V , it may be possible that $E_V(A)^2$ is not small enough. As described in the paper, in such a case one can reduce errors by averaging over many realizations of the gTPQ states because

$$\ln \overline{\langle \beta\mu V | \beta\mu V \rangle} = -V\beta j(1/\beta, \mu; V), \quad (9)$$

$$\frac{\overline{\langle \beta\mu V | \hat{A} | \beta\mu V \rangle}}{\overline{\langle \beta\mu V | \beta\mu V \rangle}} = \frac{\overline{\{\hat{A}\}'_{\beta\mu V}}}{\overline{\{\hat{1}\}'_{\beta\mu V}}} = \langle \hat{A} \rangle_{\beta\mu V}^{\text{ens}}. \quad (10)$$

Averaging over M realizations reduces the error by the factor of $1/\sqrt{M}$.

Note that averages should be taken as indicated in Eqs. (9) and (10). Otherwise, averages would not approach the correct values as shown in Supplemental Material of Ref. [2].

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[2] S. Sugiura and A. Shimizu, Phys. Rev. Lett. **111**, 010401 (2013).