

MECHANISM

To see the property of $|\psi_k\rangle$,
we take a basis set of energy eigenstates
(although we never employ such a basis
in practical calculations)

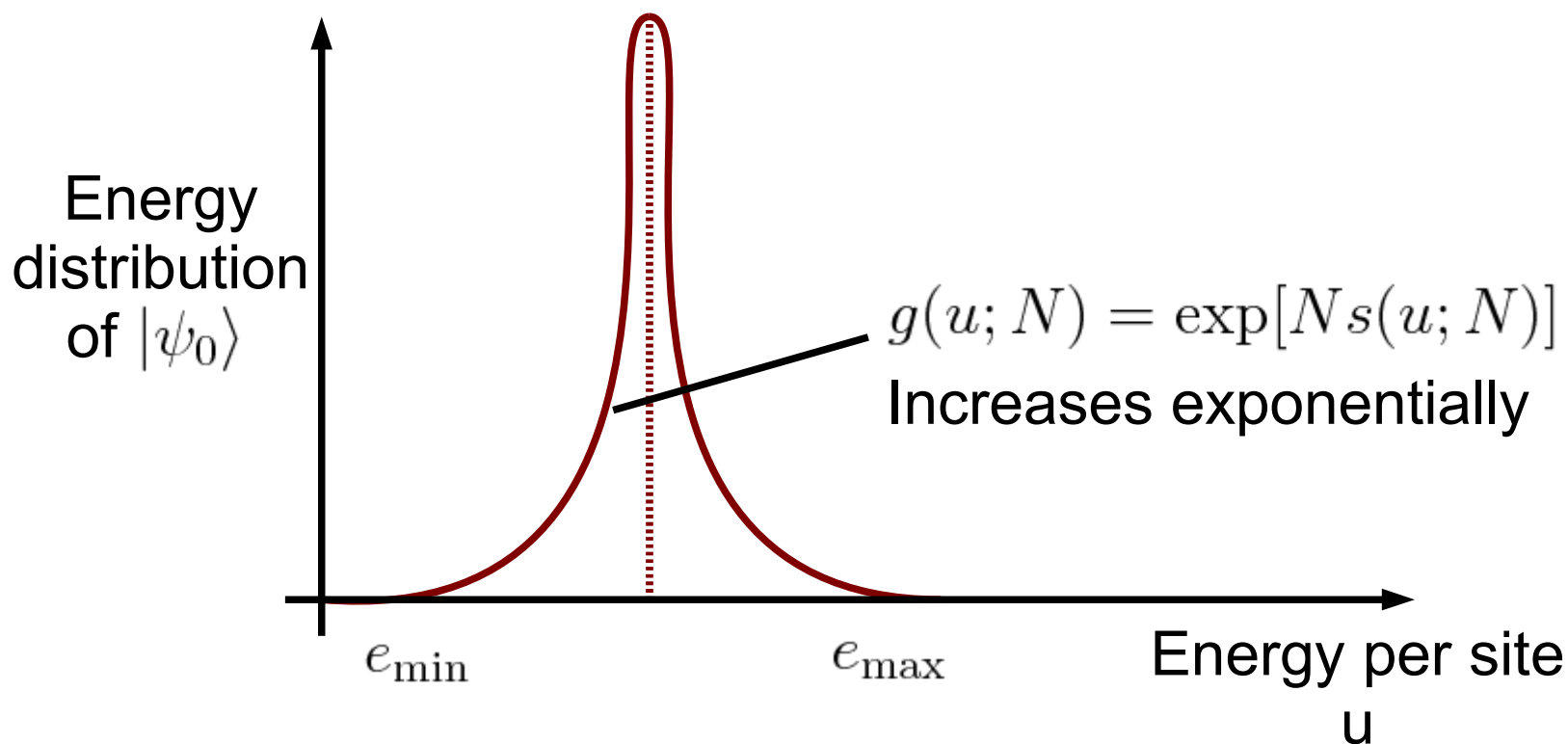
$$|\psi_0\rangle = \sum_n c_n |n\rangle \quad (\hat{h}|n\rangle = e_n |n\rangle)$$

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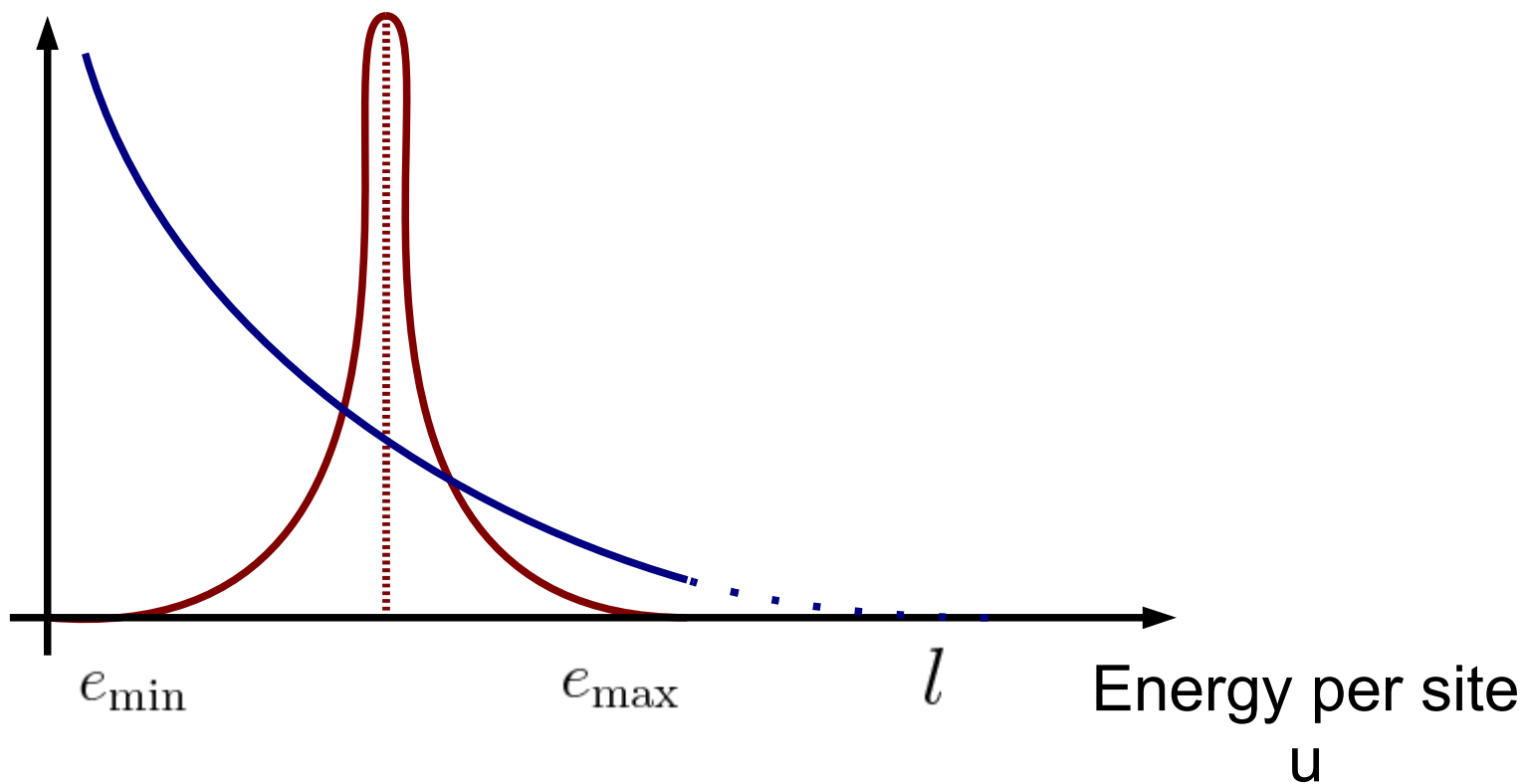
Energy distribution of $|\psi_0\rangle = \sum_n c_n |n\rangle$ ($\hat{h}|n\rangle = e_n |n\rangle$) becomes

$$\sum'_n |c_n|^2 \propto g(u; N)$$

s.t $e_n \in [u - \delta/2, u + \delta/2)$ ($g(u; N)$: density of states)



$$\begin{aligned}
 |\psi_k\rangle &\propto (l - \hat{h})^k |\psi_0\rangle \\
 &= \sum_n c_n (l - e_n)^k |n\rangle
 \end{aligned}
 \qquad (l > \|\hat{h}\|)$$

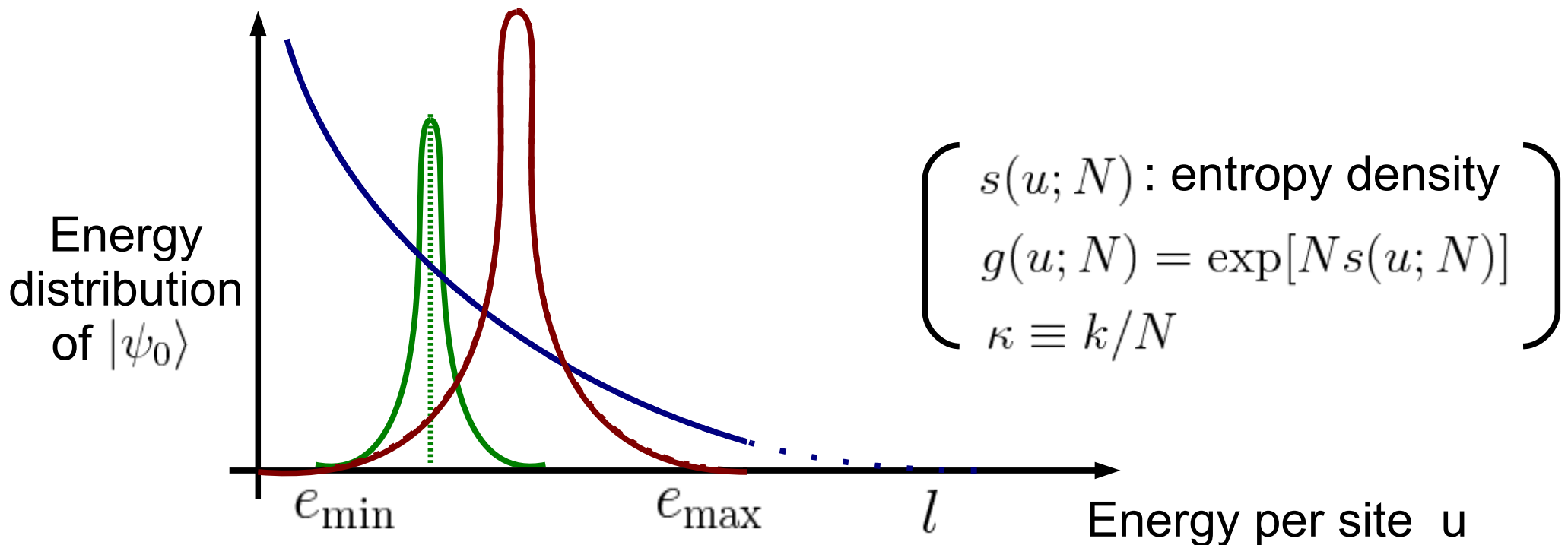


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$$\left[|\psi_0\rangle = \sum_n c_n |n\rangle \quad (\hat{h}|n\rangle = e_n |n\rangle) \right]$$

$$\sum'_n |c_n|^2 (l - e_n)^{2k} \propto g(u; N) (l - u)^{2k}$$

$$= \exp[N\{s(u; N) + 2\kappa \ln(l - u)\}]$$

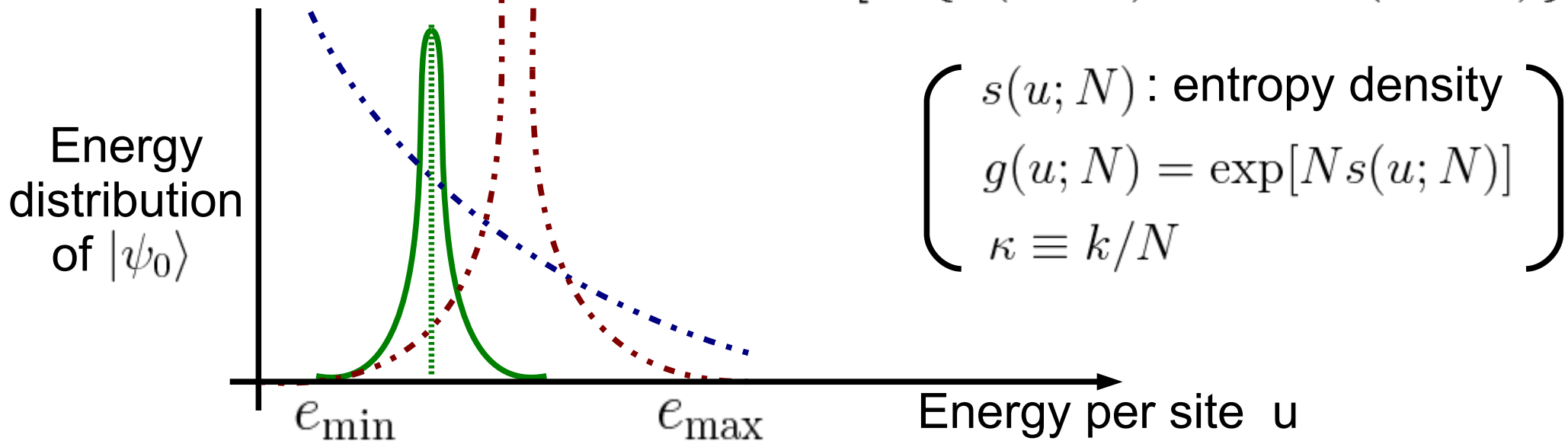


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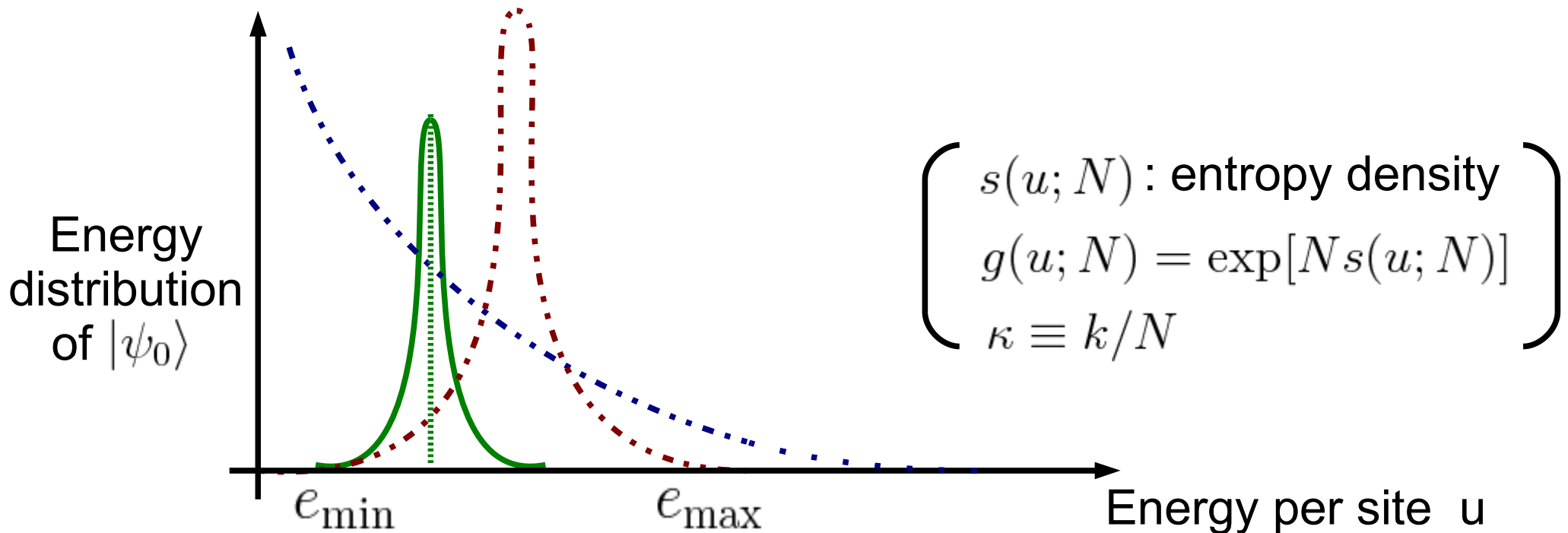
We manipulate a superposition of unknown energy eigenstates by multiplying the polynomial of Hamiltonian.

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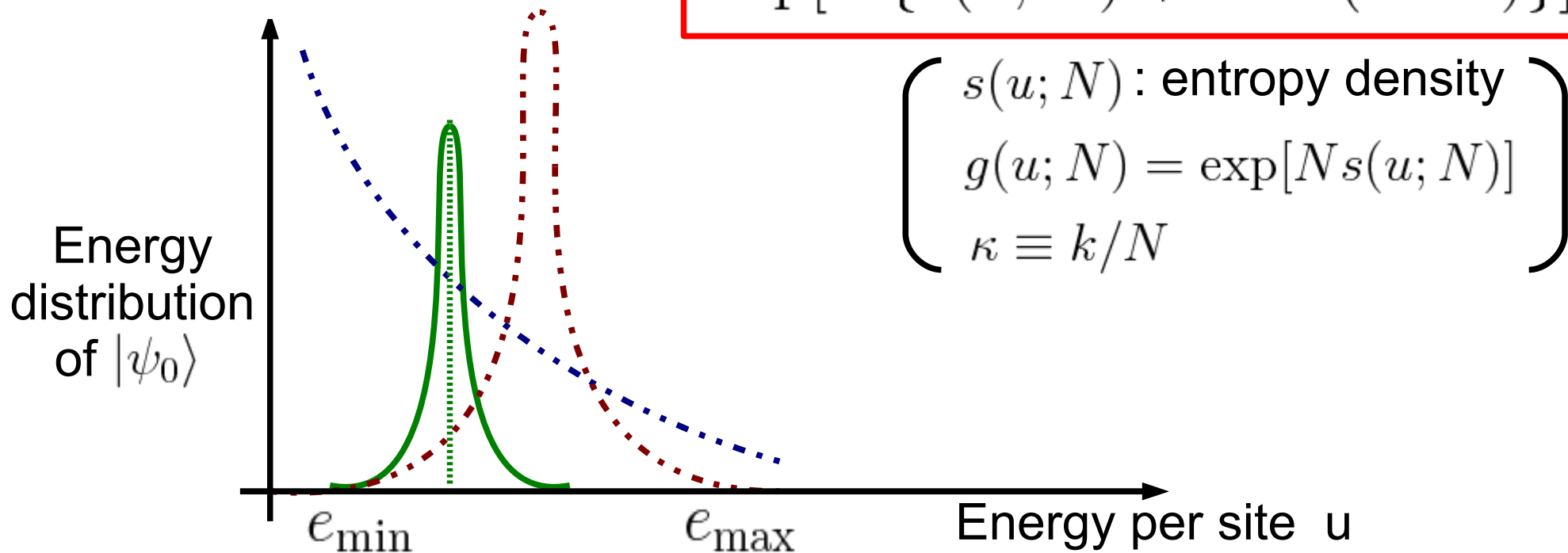
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$\xi_\kappa(u; N) \equiv s(u; N) + 2\kappa \ln(l - u)$ takes
the maximum at u_κ^* which satisfies

$$\beta(u_\kappa^*; N) = \frac{2\kappa}{(l - u_\kappa^*)}.$$

$$\left(\begin{array}{l} \beta(u; N) \equiv \frac{\partial s(u; N)}{\partial u} \\ \kappa \equiv k/N \end{array} \right)$$

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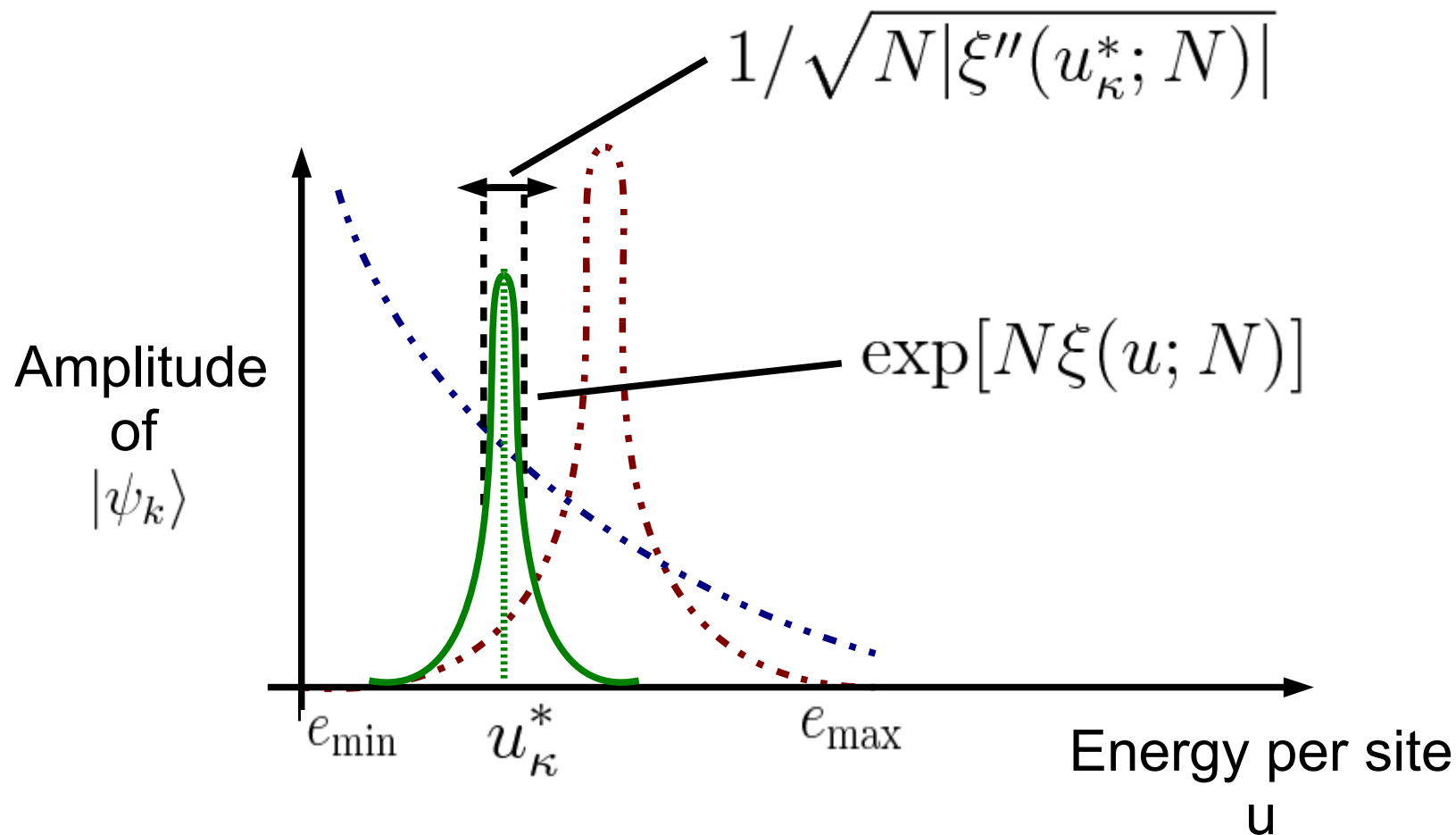
We get

$$\xi_\kappa(u; N) = \xi_\kappa(u_\kappa^*; N) - \frac{1}{2} |\xi_\kappa''| (u - u_\kappa^*)^2 + \frac{1}{6} \xi_\kappa''' (u - u_\kappa^*)^3 + \dots$$

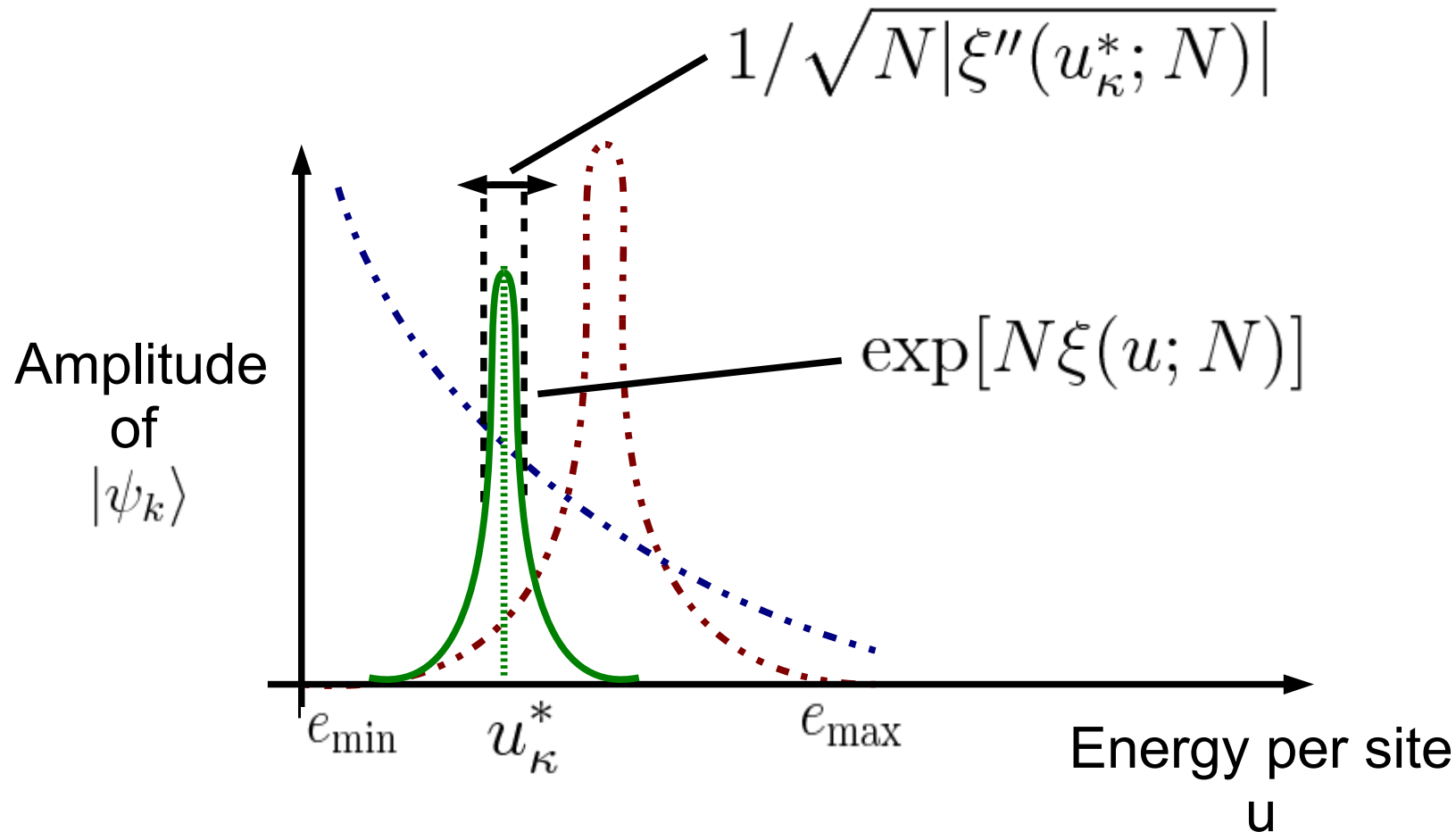
$$\left(\xi_\kappa''(u_\kappa^*; N) = \beta'(u_\kappa^*; N) - 2\kappa/(l - u_\kappa^*) \right)$$

\nearrow_0

$$\xi_{\kappa}(u; N) = \xi_{\kappa}(u_{\kappa}^*; N) - \frac{1}{2}|\xi_{\kappa}''|(u - u_{\kappa}^*)^2 + \frac{1}{6}\xi_{\kappa}'''(u - u_{\kappa}^*)^3 + \dots$$



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This means that $\langle \psi_k | \hat{A} | \psi_k \rangle$ approaches the microcanonical ensemble average of \hat{A} as $N \rightarrow \infty$.